

DEPENDENCE OF ANOMALOUS MAGNETORESISTANCE ON
TYPE OF ORIENTATION OF THE WEAK MAGNETIC
FIELD AND IRRADIATION ENERGY IN REACTOR
SHIELDING MATERIALS AT 300 K

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ABSTRACT

The longitudinal and transverse magnetoresistances (LMR and TMR) were measured of plasticized ilmenite limonite concrete sample (P.L.I.L.C). The MR was measured as a function of magnetic field (B) for different rotational angles θ (θ is the angle between the direction of both weak B and current I) at 300 K. These measurements were done again after the sample was irradiated in dry channel No 36/6 of the E-RR-I reactor. The integrated fast neutron flux > 1 Mev is 8.3×10^{17} n/cm². Different types of anomalous MR were observed especially at the two opposite directions of B before and after the sample was irradiated. Also different shapes of negative MR were found at $\pm B\theta$. From these data it can be seen that the type of anomalous MR and the intensity of negative area of MR are strongly dependent on the type of orientation of weak B, the value of opposite angle (θ) and irradiation energy. The experimental data are illustrated and the results are discussed.

INTRODUCTION

The negative magnetoresistance (nMR) which accompanied by different anomalous effects was observed in dilute magnetic alloys¹). This phenomenon was interpreted on the basis of S-d exchange interaction. Non-linear field effects in magnetic system were investigated²). They developed the statistical mechanical theory to treat non-linear effects of static magnetic fields (B) on systems of spin

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bearing particulate to liquids, and configurationally disordered (e.g. glasses) as well as rigid lattice systems.

The MR was measured in indium and $\text{Fe}_{40}\text{Ni}_{40}\text{B}_{20}$, $\text{Fe}_{30}\text{Ni}_{40}\text{Cr}_{10}\text{B}_{20}$ amorphous alloys by pulse technique in high magnetic fields³). The anomalous results indicated that the main contribution to electron conductivity of In was due to the carrier of the hole Fermi surface. They assumed that the low field behaviour of the nMR arising from magnetic effects which was masked by the positive MR arising from Lorentz force.

The aim of this work is to study the effects of both the orientation of weak B and irradiation energy (E_{irr}) on the behaviour of MR of new concrete mixer. These measurements were done under low electric and magnetic fields at $T=300\text{ K}$.

EXPERIMENTAL

Sample preparation

The concrete shield under investigation had prepared from naturally occurring Egyptian ilmenite (coarse aggregate) and limonite (fine aggregate) ores mixed with portland cement and water. A plasticizer material of 1.17 gm/m^3 density had been added to the above mixture in order to improve both mechanical and the neutron attenuation properties. Concrete mixes were designed to flow the absolute volume method recommended by the American concrete institute (A.C.I). The plasticized ilmenite limonite concrete shield (P.L.I.L.C) ingredients of 1 m^3 are shown in Table 1.

Table 1 : Ingredients of 1 m³ of PL.I.L.C shield

Material		Content
Coarse aggregate, Kg. (ilmente ore)	40 - 60 mm	722
	20 - 40 mm	482
	10 - 20 mm	482
	5 - 10 mm	241
Fine aggregate, Kg. (Limonite ore)		415
Cement content, Kg.		350
Wate, lit		287
Water cement ratio		0.82
Plasticizer material		2.25

Sample irradiation

The concrete mixer was pulverized to very fine powder and shaken very well to be homogenous. Many samples were compressed by a piston into a form of discs with equal diameter (1 cm) and equal thickness (4 mm). The intensities of the samples were measured to verify that the atoms are equally distribution in every sample ($d=2.92 \text{ gm / cm}^3$). Part of these samples were irradiated in the dry channel No 36/6 of the ET-RR-I reactor. The integrated fast neutron $> 1 \text{ Mev}$ is $8.3 \times 10^{17} \text{ n/cm}^2$. The amount of irradiation in the samples were measured outside the reactor during equal interval times (every week) until the decay of intensity of irradiation energy (E_{irr}) became normal, i.e.the samples did not radiate any amount of irradiation. These measurements were performed by using a G.M. counter.

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Measurements of magnetoresistivity

In Fig.1 terminals 1 and 2 are used for measuring the current passing the sample, while terminals 3,5 and 4,6 are used for measuring the magnetoresistivity ($\Delta\xi/\xi_0$), $\Delta\xi/\xi_0 = \xi_B - \xi_0 / \xi_0$

where ξ_B is the resistivity of the sample in the presence of B and ξ_0 is the resistivity of sample in the absence of B. All electrodes were checked to be ohmic in contact at the investigated temperature ($T = 300$ K). The measurements of LMR ($J//B$) and TMR ($J \perp B$) were done under a weak electric field ($E = 0.05$ KV / cm), where the carriers do not become hot, and at low magnetic field up to 1.5 KG to avoid the thermomagnetic effect.

RESULTS AND DISCUSSION

The dependence of LMR and TMR on the weak B for different rotational angles θ (θ is the angle between the direction of both B and I) at $T = 300$ K for PL.L.C sample are shown in Figs.2 and 3. The general shapes of B-Mr relationships at all θ are quadratic dependences on B, i.e. $\Delta\xi/\xi_0 \propto B^2$. Where every parabolic curve is partially negative at low B and becomes positive at relatively high B. Figs.1 and 2 show different types of anomalous MR. The appearance of negative MR (nMR) at relatively high temperature means that the nMR is not only low temperature dependent as is well-known up to 50K^{4, 5}), but also high temperature dependent phenomenon. In our case the nMR was appeared due to the absence of symmetry axes in concrete sample.

According to these measurements, it is interesting to analyse this experimental data which can be characterized with the following fundamental features :

1 - The rate of intensity variation of nMR in LMR is four times larger than those of TMR, but the variation rate of positive LMR is five times bigger than that observed in positive TMR during the whole scale of B at all θ .

2 - The negative depth of LMR at all θ are more deeper than at TMR, but in the TMR the negative areas have boat shapes, i.e. the negative saturation cases were appeared with different degrees. The saturation case means that the nMR is B independent,

$$(-L \Delta\xi / \xi_0 B)_\theta \neq (-T \Delta\xi / \xi_0 B)_\theta \quad \text{for } T = 300 \text{ K.}$$

3 - From Figs. 2 and 3 it can be seen that the anomalous behaviours of MR are observed when the direction of B was reversed at any θ . The nMR at $\theta = 0^\circ$ is larger than that at $\theta = 180^\circ$ in Fig. 2. Also in Fig. 3 it was found that the anomalous effects of nMR at all θ except $\theta = 90^\circ$ are nearly closed from each other in magnetic field range 0.5-0.8 KG. This shows that the value θ and the type of orientation of weak B are strongly effective on the quantitative and qualitative behaviours of MR; $(L \Delta\xi / \xi_0 B \pm \theta), (T \Delta\xi / \xi_0 B \pm \theta) \neq \text{const}$ for $T = 300 \text{ K.}$

4 - It is well known that the measurements of MR as a function of $\pm B$ under low electric and magnetic field are isotropic⁶). But in our measurements the effect of $\pm B_\theta$ on MR is asymmetric. Where the degrees of stability or instability of nMR are controlled by a certain compensating relation between the directions of $\pm\theta$ and the type of orientation of B.

The PL.L.C sample contains three basic compounds, Fe_2O_3 (31.2%), TiO_2 (26.5%) and FeO (19.1%) by weight and many traces of elements as appeared from Tefa analysis (Table.2). The atoms of these compounds have a good ability to

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change their arrangements when exposed to the external magnetic field, due to the high magnetic moment of these elements (the total magnetic moments (in units U_B) of Fe = 6, Ti = 4 /3 and Ca = 0⁷).

From the above mentioned it was found that the different compositions of PL.I.L.C sample and the high value of magnetic moment of iron which were created a numerous asymmetry axes. This system is responsible for the appearance of different types of anomalous effects of MR.

Table. 2 : Chemical compositions of PL.I.L.C. sample of density 2.95 / cm³

No	Material	wt. %	Density	Mol. wt.
1	Fe ₂ O ₃	31.2	5.24	159.69
2	TiO ₂	26.5	5.248	79.9
3	FeO	19.1	5.7	71.85
4	CaO	7.89	3.32	56.08
5	SiO ₂	6.4	2.64	60.08
6	H ₂ O	6.29	1	1
7	Al ₂ O ₃	1.47	3.97	101.96
8	Plasticizer material	0.7	1.17	
9	MgO	0.25	3.58	40.31
10	So3	0.2	2.29	80.6
Total		100 %	2.95	

The accurate data of the analysis of X-ray diffraction of PL.I.L.C sample before and after irradiation.

The sample was irradiated in the dry channel No. 36/6 of the ET-RR-I reactor. The integrated fast neutron > 1 Mev is 8.3×10^{-17} n/cm².

The X-ray diffraction patterns were automatically recorded by using a modern Philips diffractometer. The general operating conditions were as following :

Copper K- alpha radiation, 38 KV, 16 mA, slits : 1° ; 0.2 mm; 1° proportional counter and scanning velocity $2^\circ(2\theta)$ / min.

The two diffraction patterns of PL.I.L.C samples before and after irradiation (B_{irr} and A_{irr}) are shown in Fig.4. It was observed that the two patterns B and A have similar behaviours at the same angles (2θ) before and after the sample was irradiated but in different intensities. On the other hand the intensity of the peaks of pattern B is much sharper than that in pattern A at the same 2θ . From Fig.4, it can be seen that the amount of proper amorphous part does not exceed 56% of the total compositions of the sample.

From Table.3. it was found that the spectrum lines of some compounds are the same B_{irr} and A_{irr} at certain 2θ . This shows that at these angles 2θ , the E_{irr} has no effect on these compounds and may be called resistive angle $(2\theta)_R$. Also the spectrum line of Fe_2O_3 at $2\theta = 33.6$ is 100% B_{irr} and A_{irr} , i.e. this value is identical with the value in the identification card. This means that the E_{irr} at this angle may be suffer from completely reflection by the crystalline planes of these atoms. On the other hand different similarity in the values at different 2θ may be due to that some parts of E_{irr} are asymmetry scattering and the remain parts may be

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reflected according to the value of 2θ . Moreover may be the E_{iRR} at these angles has very weak effect on the atoms. So the absorption and emission processes are isothermic. This means that the E_{iRR} was emitted or absorbed by the atomic planes in equal rates.

As may seen from Table 3 at definit 2θ , the intensity of spectrum lines A_{iRR} are many times weaker than B_{iRR} . This indicates that the E_{iRR} was consumed completely to convert the crystalline states to amorphous states at these angles, so these angles may be called inversion angles $(2\theta)_I$.

Table. 3 : Intensities of resistive angle $(2\theta)_R$ and inversion angle $(2\theta)_I$ of groups $(Fe_2O_3.H_2O)$, $Fe_2O_3TiO_2$ ($FeTiO_5$) and $FeTiO_3$ ($FeOTO_2$). These groups have a symbols I, II, and III respectively.

Groups	$(2\theta)_R$	$(I/I_0)_B$	$(I/I_0)_A$	$(2\theta)_I$	$(I/I_0)_B$	$(I/I_0)_A$
I	33.6	100	100	19	26	4
	41.3	30	30	21.9	33	7
				41.9	38	9
				48.2	30	7
II	36	48	48	32.5	32	8
	41.3	30	30	48.2	30	7
			60.2	32	7	
III	23.9	12	12	32.5	32	8
	41.3	30	30	48.2	30	7
	64.5	25	25			

The dependences of LMR and TMR on weak B for different θ after the sample was irradiated (A_{irr}) at $T = 300K$ are shown in Figs.5 and 6. Many types of anomalous MR were observed, Which are in contrast to the behaviours shown in Figs.2 and3 before irradiation (B_{irr}). So the comparison between the anomalous MR B_{irr} and A_{irr} can be summarized as the following.

1 - The LMR B_{irr} and A_{irr} are parabolic dependent on B^2 at all opposite angles θ , but the relative values of LMR A_{irr} are many times smaller than B_{irr} . Also the general shapes of nMR B_{irr} at all θ are more clearness than A_{irr} . Moreover there are irregular distributions among the MR curves at different θ B_{irr} and A_{irr} . This means that the E_{irr} is not only destroyed the lattice structures of crystalloid atoms but also strongly effective on the behaviour of MR.

$$(L\Delta\xi / \xi_0 B(\theta)) B_{irr} \neq (L\Delta\xi / \xi_0 B(\theta)) A_{irr} \text{ for } T = 300 \text{ K.}$$

2 - The nMR were found in magnetic field ranged 0.1-0.6 KG at all θ A_{irr} and at $B=0.1-1.6KG$ B_{irr} . This shows that the E_{irr} may be minimize the areas of MR curves due to the increasing of resistivity as a result of phase transition;

$$(nL \Delta\xi / \xi_0 B(\theta)) B_{irr} \gg (nL\Delta\xi / \xi_0 B(\theta)) A_{irr} \text{ for } T = 300K.$$

3 - The anomalous behaviours of LMR at $B \pm \theta$ B_{irr} are more clearness than A_{irr} . where the negative curves of MR A_{irr} for different θ are nearer from each other than B_{irr} . due to the decreasing rates of crystalloid states. Thus the E_{irr} is considered destructive force, i.e. the asymmetry directions of atoms are increasing A_{irr} ;

$$(L\Delta\xi / \xi_0 B(\theta)) B_{irr} \neq (L\Delta\xi / \xi_0 B(\theta + 180^\circ)) A_{irr} \text{ for } T = 300K$$

By the same manner, the above steps of comparison between the characters of

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LMR B_{irr} and A_{irr} are repeated themselves in TMR (Figs.3 and 6) but in different rates. The saturation cases of TMR at all θ were disappeared after the sample was irradiated. This indicates that absence of independence of MR on B A_{irr} is due to the random distributions of atoms of different compounds in PL.L.C. sample;

$$(T \Delta\xi / \xi_0 B_{\pm\theta}) B_{irr} \neq (T \Delta\xi_0 B_{\pm\theta}) A_{irr} \text{ for } T = 300 \text{ K.}$$

The effect of E_{irr} on the variation of anomalous behaviours of LMR and TMR is very large. The relative values of positive LMR (Fig.5) are smaller than TMR at all θ (Fig. 6), a contrast occurs in the normal case. On the other hand the nLMR at different θ are many times larger than nTMR. The two anomalous effects on the positive and negative MR may be attributed to the non-compensating relations among the E_{irr} on the type of orientation of weak B, (E_{irr}, θ) , (E_{irr}, B) and $(E_{irr}, \sigma B)$. These measurements indicate that the anomalous behaviours of MR are not only low B, θ , T, type of orientation of B^{8-21} dependent, but also E_{irr} dependent phenomenon;

$$(+ve L \Delta\xi / \xi_0 B_{\pm\theta}) A_{irr} < (+ve T \Delta\xi / \xi_0 B_{\pm\theta}) A_{irr};$$

$$(-ve L \Delta\xi / \xi_0 B_{\pm\theta}) A_{irr} > (-ve T \Delta\xi / \xi_0 B_{\pm\theta}) A_{irr};$$

$$(L \Delta\xi / \xi_0 B_{\pm\theta}) A_{irr} \neq (T \Delta\xi / \xi_0 B_{\pm\theta}) A_{irr}; \# \text{ constant for } T = 300 \text{ K.}$$

To calculate the two contributions of MR, the measured data of MR could be written in the power series of B as the following :

$$(\Delta\xi / \xi_0 (B, \theta))_{T=300 \text{ K}} = a_0 B^0 + a_1 B^1 + a_2 B^2 + \dots$$

where the first order term $(\Delta\xi / \xi_0 \sim B^1)$ describes the non-parabolic term of MR

and the second order term $\Delta\xi / \xi_0 \sim B^2$ represents the conventional positive parabolic term of MR.

The constants a_0 , a_1 and a_2 are dependent on the geometry of the sample (orientation), type of the material and temperature.

The above empirical formula was fitted by the measured LMR and TMR before and after the sample was irradiated for calculating the two components of MR as shown in Figs. 7 - 14.

The linear terms of MR are exhibited many types of anomalous effects B_{irr} and A_{irr} . The irregular distributions of non-parabolic terms of LMR and TMR at $\pm B_\theta$ as shown in Figs.7 and 8 are strongly dependent on the type of orientation of weak B and θ , i.e. at $B = 0.1$ KG (Fig.8) there are no anomalous effects, except at $\theta = 90^\circ$. These behaviours of linear terms of TMR are completely different from in LMR (Fig. 7), i.e. at $B = 0.1$ KG the anomalous behaviours of all components are appeared except at $\theta = 90^\circ$ and 180° ;

$$\left| T a_1 B \right|_{\pm\theta} \neq \text{cons}; \quad \text{for } B = 0.1 \text{ KG, except}$$

$$\theta = 90^\circ$$

$$\left| L a_1 B \right|_{\pm\theta} \neq \left| T a_1 B \right|_{\pm\theta}. \quad \text{for } T = 300\text{k}$$

After the sample was irradiated, it was found that the anomalous effects of LMR and TMR (Figs.9 and 10) are in contrast to the behaviour shown in Figs.7 and 8. In Fig.9 it was noticed that the anomalous gap of $L a_1 B^1$ at $\theta = (45^\circ, 225^\circ)$ is larger than that at $\theta = (0^\circ, 180^\circ)$. Also there is no consider gap observed at $\theta = 0^\circ$ and 270° . In Fig.10 the three linear terms of TMR at $\theta = 45^\circ, 90^\circ$ and 225° are

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closed from each other and at low B they not exhibited any type of anomalous effects. The relative calculated values of L_{a_1B} and T_{a_1B} along the whole scal of weak B are nearly the same, but the effect of $\pm\theta$ on a_1B is different,

$$|L_{a_1B}(\theta)|_{A_{irr}} \# |L_{a_1B}(\theta+180^\circ)|_{A_{irr}};$$

$$|T_{a_1B}|_{\theta = 45^\circ, 90^\circ, 225^\circ} \# \text{Cons for } B = 0.3 \text{ KG};$$

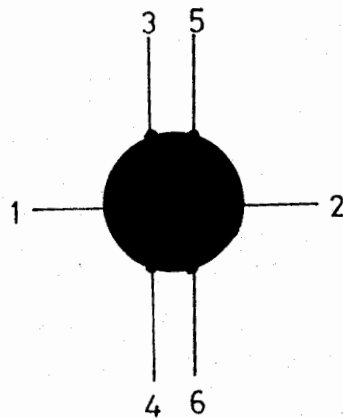
$$|L_{a_1B}, T_{a_1B}(\pm\theta)| \# L_{a_1B}, T_{a_1B}(\pm\theta) |_{A_{irr}} \text{ for } T = 300 \text{ K.}$$

In Figs. 11 - 14, the calculated normal positive components of LMR and TMR B_{irr} and A_{irr} for all θ are linear dependent on B^2 as will knwn. The anomalous behaviours of these terms were observed at two opposite directions of B, but in different rates in both quantitative and qualit behaviours. The intensity of asymmetry relations of parabolic components are controlled by the E_{irr} , the values of $\pm\theta$ and weak B,

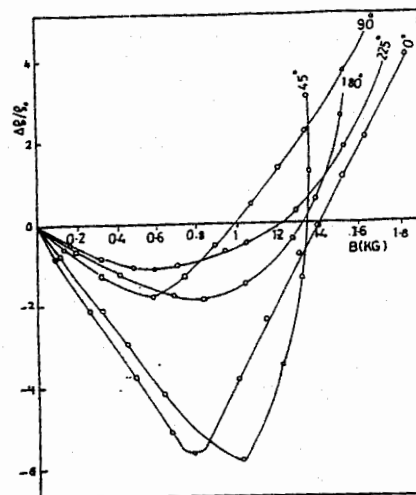
$$|a_2B^2|_{\theta} \# |a_2B^2|_{(\theta+180^\circ)} \text{ for LMR and TMR } B_{irr}$$

and A_{irr} at $T=300 \text{ K.}$

To examin the concept of the experimental MR as a function of power series of B, which is equal is to the summation of the two calculated terms. One of them is negative, odd power of B, and the other is positive, even power, of B, so the relation between the two terms with experimental curve of MR and B at some different θB_{irr} and A_{irr} are shown in Figs. 15 - 18. From these Figs, it can be seen that the above concept of MR summation is valid and remain as it is after the sample was irradiated. The change of positions and relative values of the two components are due to the variation of θ, B and the effect of E_{irr} on the asymmetry

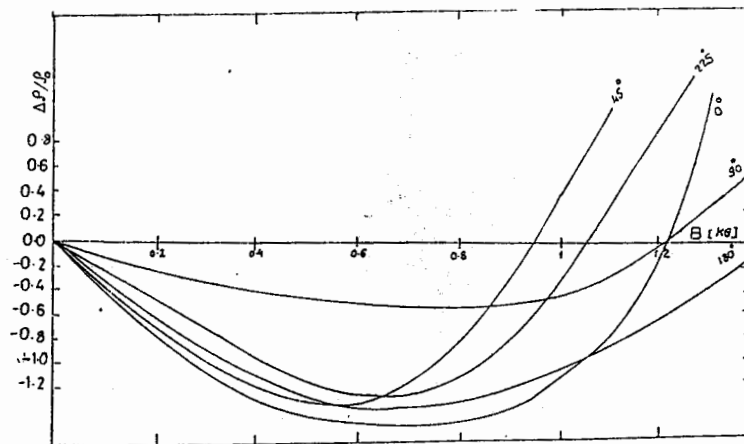


(Fig. 1): The shape of PL.I.L.C. sample (disc form).

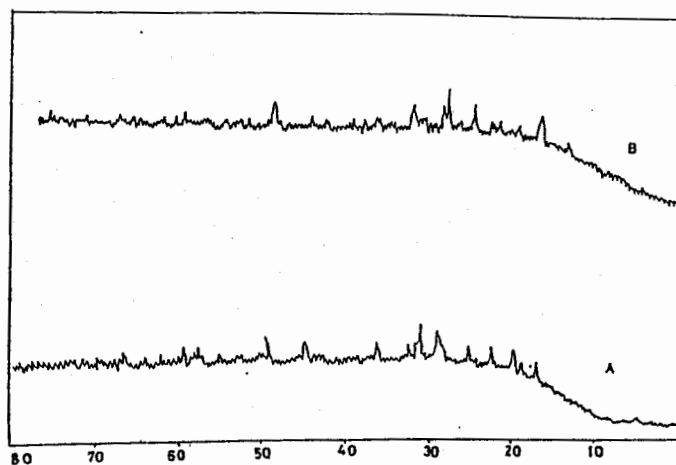


(Fig. 2): The measured longitudinal magnetoresistivity ($L\Delta\xi/\xi_0$) as a function of magnetic field (B) for different rotational angles (θ) at $T=300$ K.

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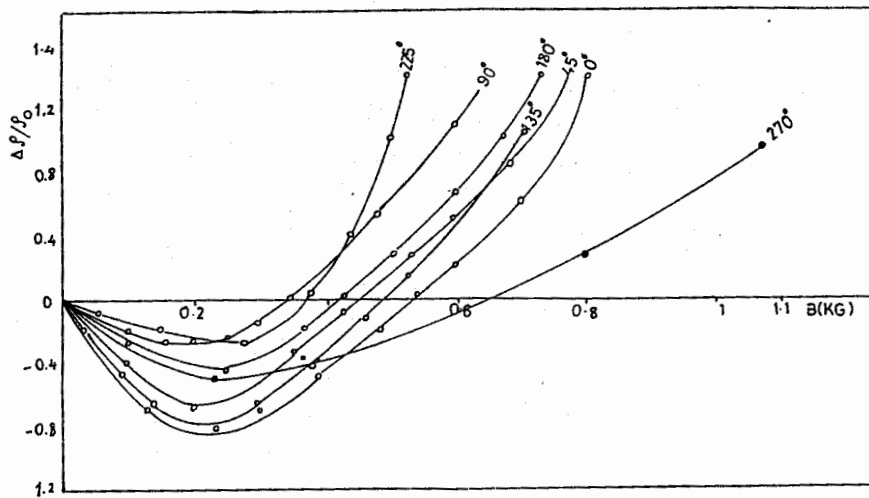


(Fig. 3) : The measured transverse D_r / r_0 ($T \Delta \xi / \xi_0$) as a function of B for different θ at $t = 300$ K.

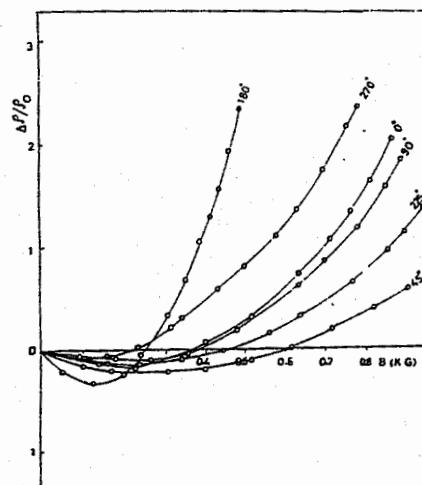


(Fig. 4) : X- ray diffraction pattern of PL.I.L.C. sample before and after the sample was irradiated at $T= 300$ K.

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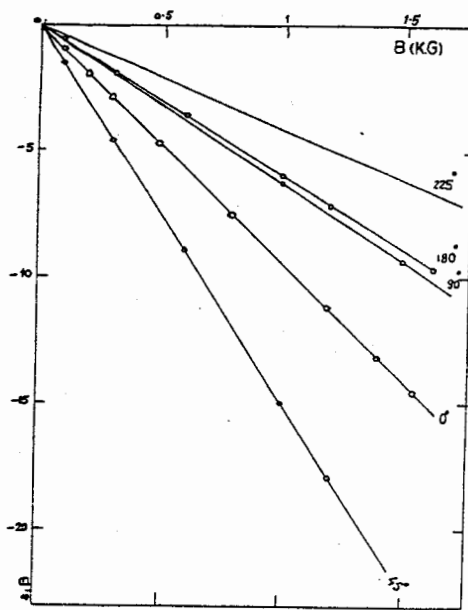


(Fig. 5) : Dependence of $L \Delta\xi / \xi_0$ on B at $T = 300$ K for different θ , when the sample was irradiated.

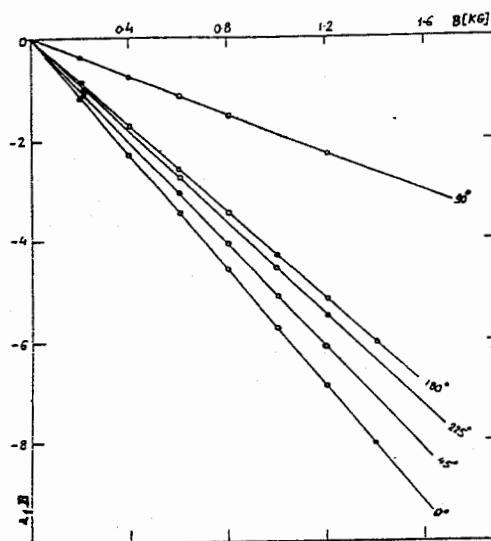


(Fig. 6) : Dependence of $T \Delta\xi / \xi_0$ on B at $T = 300$ K for different θ , when the sample was irradiated.

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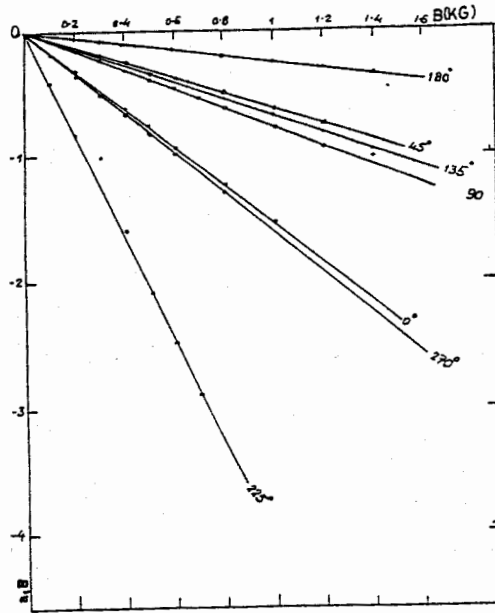


(Fig. 7): The calculated odd components ($a_1 B$) of $L \Delta \xi / \xi_0$ versus B for different θ at $T = 300$ K.

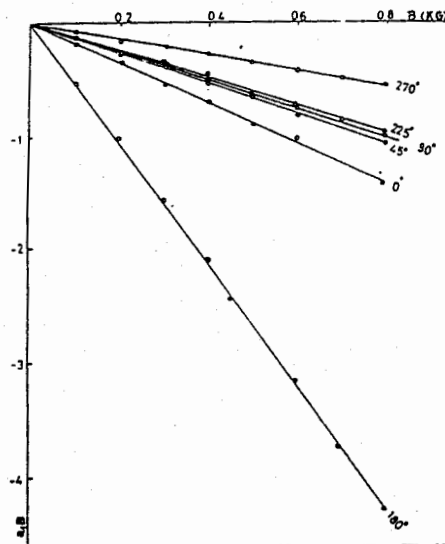


(Fig. 8): The calculated $a_1 B$ of $T \Delta \xi / \xi_0$ versus B for different θ at $T = 300$ K.

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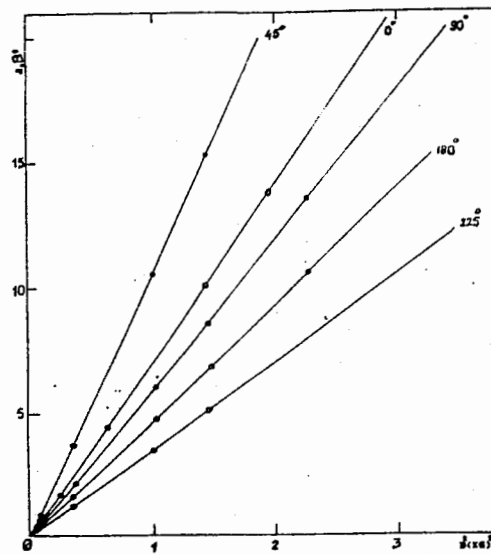


(Fig. 9) : The Calculated a_1B of $L \Delta\xi / \xi_0$ versus B for different θ at $T = 300$ K, when the sample was irradiated.

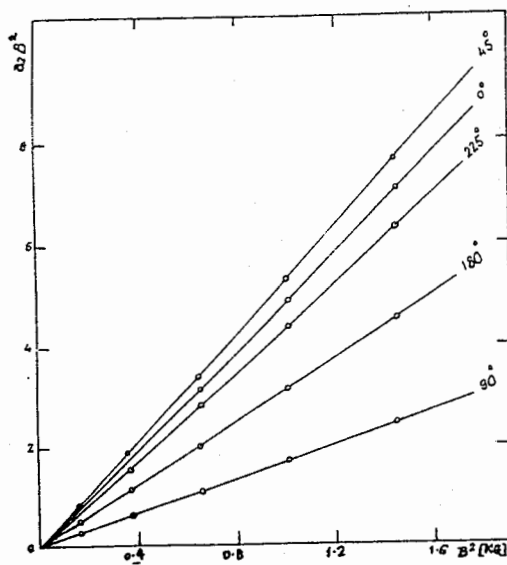


(Fig. 10) : The Calculated a_1B of $T \Delta\xi / \xi_0$ versus B for different θ at $T = 300$ K, when the sample was irradiated.

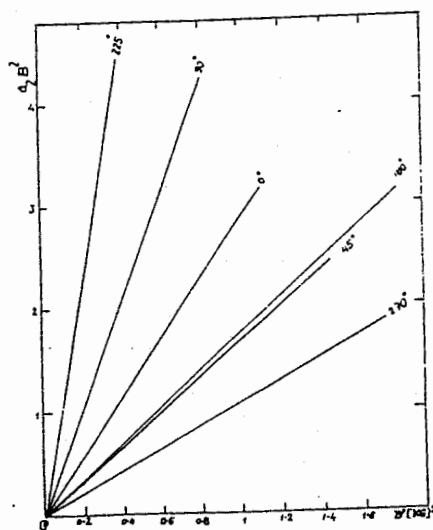
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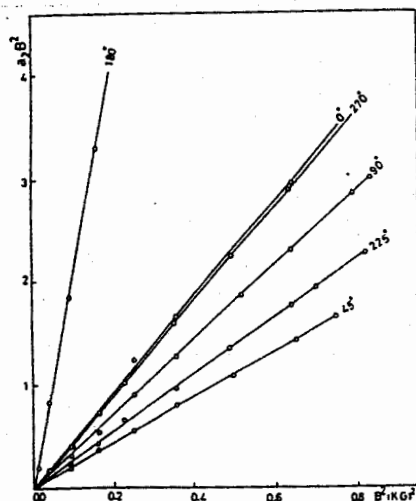
(Fig. 11) : The Calculated even components $a_2 B^2$ versus B at several θ for $L \Delta \xi / \xi_0$ at $T = 300$ K.



(Fig. 12) : The Calculated $a_2 B^2$ versus B at several θ for $T \Delta \xi / \xi_0$ at $T = 300$ K.

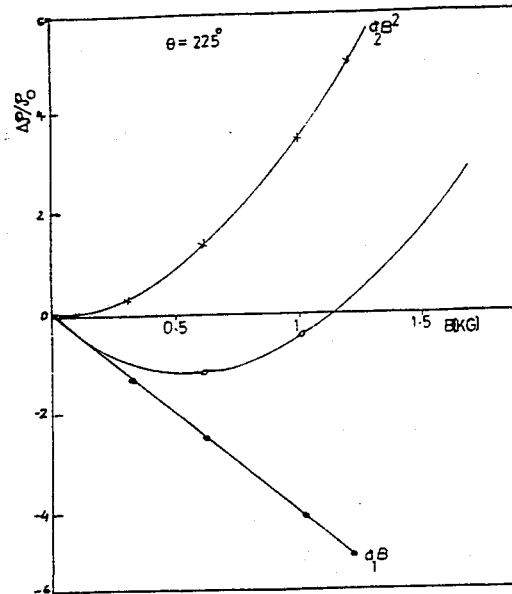


(Fig. 13) : The Calculated $a_2 B^2$ versus B at several θ for $L \Delta \xi / \xi_0$ at $T = 300$ K, when the sample was irradiated.

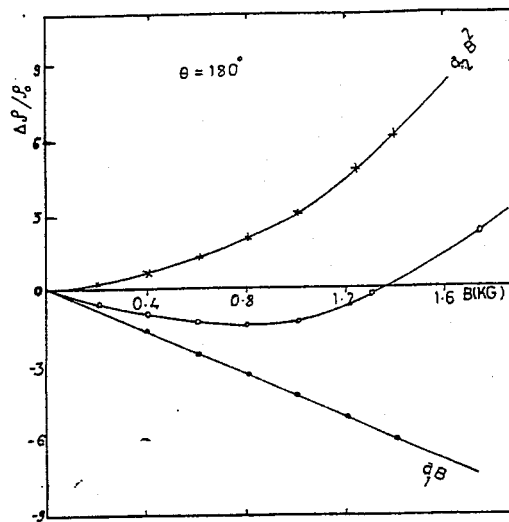


(Fig. 14) : The Calculated $a_2 B^2$ at several θ for $T \Delta \xi / \xi_0$ at $T = 300$ K, when the sample was irradiated.

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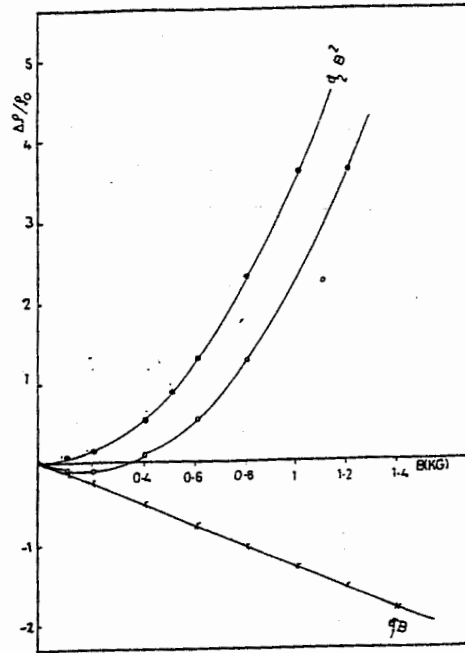


(Fig. 15): Dependence of the measured $L\Delta\xi/\xi_0$ on B for $\theta = 225^\circ$ at $T = 300$ K. The calculated anomalous a_1B and normal a_2B^2 are represented by symbols (x) and -o- respectively.

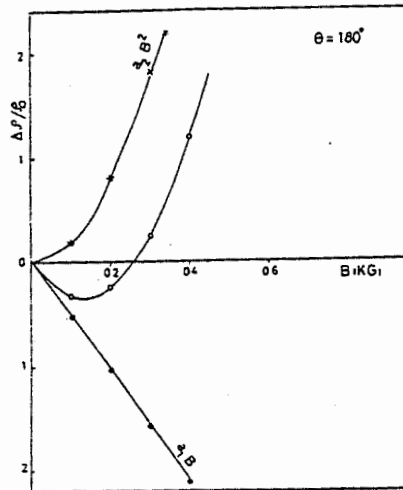


(Fig. 16): Dependence of the measured $T\Delta\xi/\xi_0$ (-o-) on B for $\theta = 180^\circ$ at $T = 300$ K. The calculated anomalous a_1B and normal a_2B^2 are represented by symbols (-o-) and (x) respectively.

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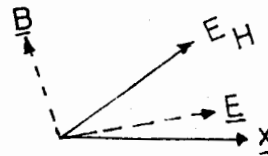


(Fig. 17) : Dependence of the measured $L\Delta\xi/\xi_0$ (-o-) on the B for $\theta = 90^\circ$ at $T = 300$ K. The calculated anomalous a_1B and normal a_2B^2 are represented by symbols(x)and (-o-) respectively after the sample was irradiated.

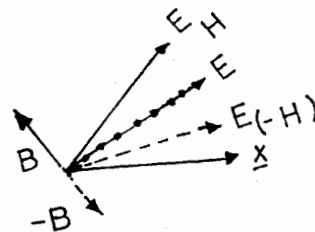


(Fig. 18) : Dependence of the measured $T\Delta\xi/\xi_0$ (-o-) on B for $\theta = 180^\circ$ at $T = 300$ K. The calculated anomalous a_1B and normal a_2B^2 are represented by symbols(-o-) and (x) respectively after the sample was irradiated.

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(Fig. 19) : Direction of induced Hall field (E_H) as a result of application of weak electric (E) and magnetic (B) fields, where X is the preferred direction.



(Fig. 20) : Illustrated asymmetry relation between the induced E_H at the two opposite directions of B .

axes of compounds of PL.I.L.C sample ;

$$\langle \Delta \xi / \xi_0 B(\theta) \rangle = \sum \# \text{La}_1\text{B}, \text{Ta}_1\text{B} (\pm\theta) \mid A_{\text{irr}} \text{ for } T = 300 \text{ K.}$$

The ratio of errors in the above expression are different according to the type of orientation and the value of θ , B_{irr} and A_{irr} . In general the errors of the above calculations were ranged from decimal fraction or hundredth.

The above results may be attributed to the absence of what is called symmetry axes in the PL.I.L.C sample. Any aggregate is characterized by a small parts have a crystalline behaviours and may called quasi- crystals or crystalloids. This system is said to have a preferred orientation, or texture, which each grain in the polycrystalline part normally has a crystallographic orientation, which is different from that of its neighbours. Considered as a whole, the orientations of all grains may be randomly distributed in relation to the same selected fram of reference, or may tend to clustor, to a greater or degree about some particular orientation or orientations.

The main effect of E_{irr} on the preferred orientations or texture is restrained between the distructibility of texture in different rates and vary the positions of preferred orientations. This means that befor and after the sample was irradiated there are asymmetric axes in disordered and ordered states, i.e. the direction of induced Hall field $E_{\text{H}}B_{\text{irr}}$ or A_{irr} is not perpendicular to the directions of both electric and magnetic fields as shown in Fig.19 The absence of symmetry directions in the PL.I.L.C sample are caused in created a certain type of electric potential gradient where the intensity of this field is strongly dependent in the θ and E_{irr} . So according to this assumption the direction of E_{H} at the two opposite directions of B are

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asymmetrical as shown in Fig. 20.

The different types of iron oxide groups in the PL.I.L.C sample in the presence of B are induced a magnetic areas on the two sides of surfaces sample. These magnetic areas are dependent on the type of orientation of B and the intensity of interaction between the internal magnetic and external B. After the sample was irradiated the rate of induced magnetic areas are changed. All the above reasons are responsible for the variations of dynamic behaviours of different carriers such as effective mass (m^*), conductivity (δ), mobility (μ) and...etc;

$$\Delta\zeta/\xi_0 = \frac{\sigma(0) - \sigma(B)}{\sigma_B} = \frac{\mu(0) - \mu(B)}{\mu(B)}$$

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