## **INVESTIGATION OF NEUTRON DOSES IN WATER**

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#### ABSTRACT

Thermal, intermediate and fast neutron fluxes from compact D-T neutron generator have been measured inside a tissue equivalent medium(water). The method is based on the measuring of induced activities produced by the activation foil technique. The fluence dose conversion factors as well as related quality factors were used to determine the dose equivalent. The contribution of thermal and intermediate neutron doses is less than 1% of the total dose.

### **INTRODUCTION**

The use of neutrons in therapy requires accurate physical dosimetry. It is very important to specify in details the spatial distribution of absorbed dose rate in a tissue equivalent medium [1-3]. D-T neutron distributions were determined inside the media by using the activation foil technique [4-5]. If one has a measure of neutron flux then the dose rates could be calculated using the ICRP data [6].

The question of neutron spectrum dependence must be raised for calculated as well as measured data. Size of detector has a great importance as each measure of the dose at a point involves integration of energy deposition over the region of the detector [7]. The use of activation foils as a neutron detectors [8] shows by their small thickness good properties close to the neutron source.

The aim of this work is to investigate the flux and dose distribution for D-T neutrons inside a water phantom with dimensions (30x30x20) cm<sup>3</sup>.

#### **EXPERIMENTAL**

The measurements were carried out with D-T compact neutron source of Phillips type PW5310, yield  $1.1 \times 10^8$  n.s<sup>-1</sup> and 14.5 MeV neutron energy. The spectrum of generator was examined by activation foil technique [5]. Tab. 1 gives the characteristics of the used foils and the considered reactions which used to determine either fast or thermal neutron fluxes as well as the convenient times of irradiations.

In addition cadmium foils were used to cut off thermal neutrons from diffused spectrum. Foils were supplied by BDH England and Reactor Experiments Inc. San Carlos, California, U S A.

| Foil | Radius<br>(cm) | Thickness<br>(mm) | Reaction of interest                                   | Threshold<br>energy | T <sub>1/2</sub> | Ti        |
|------|----------------|-------------------|--|---------------------|------------------|-----------|
| Zn   | 1.30           | 1.83              | 64Zn(n,2n) $63$ Zn                                     | 12.10               | 38.0m            | 40.0m     |
| Cu   | 1.40           | 0.10              | $63_{Cu(n,2n)}62_{Cu}$                                 | 11.00               | 9.8m             | 10.0m     |
| Al   | 1.84           | 1.00              | $\frac{27_{Al(n,\alpha)}^{24}Na}{27_{Al(n,p)}^{27}Mg}$ | 3.25<br>1.90        | 15.0h<br>9.4m    | 6h<br>10m |
| In   | 1.37           | 1.00              | $\frac{115_{In(n,n')}115m}{115_{In(n,\gamma)}116m}$ In | 0.34                | 4.5h<br>54.2m    | 2h<br>2h  |

Table 1: The characteristics of the used foils.

$$\mathbf{A} = \mathbf{C} \int_{\mathbf{E}_{t}}^{\mathbf{E}_{m}} \sigma(\mathbf{E}) \, \phi(\mathbf{E}) \, \mathbf{d}\mathbf{E}.....$$

The produced apparent activities of the radioproducts (table 1) were measured by NaI(Tl) scintillation crystal 3" x3" connected to a multichannel analyzer Canberra series 40. Fig. 1 gives the efficiency of the used spectrometer. Proper corrections were applied to the apparent activity in order to get the actual activity A. The differential flux in each energy region was determined by applying the successive stepwise approximation method [5], where A is related to the differential flux  $\phi$  (E) by the relation:

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(1)

The elemental mass of the detector, the relative abundance, the atomic weight, Avogadro's number and corrections for activation and decay time are coupled together as constant C.  $\sigma$  (E) is the energy integral of the activation cross section between  $E_t$  and  $E_m$  where  $E_t$  is the threshold activation energy of the detector and  $E_m$  is the cut-off neutron energy.  $\sigma$  (E) could be determined from the available literatures [9]. Normalization corrections had to be applied to all activations using a neutron monitor [5].

After deduction of the fast neutron flux for each energy region from 14.5 MeV to 0.34 MeV, the absorbed dose D(E) is determined by using ICRP conversion factors and calculated by using the equation(1).

$$\mathbf{D}(\mathbf{E}) = \phi(\mathbf{E}) \, \tilde{\mathbf{E}} (\bar{\sigma}_{ai} \, \mathbf{n}_i + \mathbf{F}_i \, \bar{\sigma}_{si} \, \mathbf{n}_i) \tag{2}$$

where:

 $\phi$  (E) is the distribution of the neutron flux with respect to energy,

**E** is the average neutron energy in each region,

 $\sigma_{ai}, \sigma_{si}$  are the average absorption and scattering cross sections of the <u>ith</u> element in each energy

region,

n; is the number of atoms for unit mass of ith element,

F<sub>i</sub> is the average fractional energy transfer to the scattered atom.

As one has calculated the absorbed doses then it is easy to find the equivalent dose rates according to the following equation [1]:

$$\mathbf{DE} = \mathbf{D} (\mathbf{E}) \mathbf{x} \mathbf{QF}$$
(3)

where DE is the equivalent dose rate and QF is the average quality factor in each energy region.

The accumulation of the scattered neutrons at deep depths within the phantom is yielding the build-up factor  $B_D(E,r)$  which is defined as:

$$\mathbf{B}_{\mathbf{D}}(\mathbf{E},\mathbf{r}) = \left(\frac{\mathbf{D}(\mathbf{E},\mathbf{r})}{\mathbf{D}(\mathbf{E},\mathbf{o})}\right) \mathbf{e}^{-\Sigma_{\iota}\mathbf{r}}$$
(4)

### **RESULTS AND DISCUSSION**

Figure(2) represents the integrated fast neutron flux overall different energy groups, while figure(3) represents the variation of intermediate and thermal neutron fluxes with the distance inside the phantom. It is clear that the integrated fast neutron flux (14.5-0.34 MeV) decreases gradually with distance which is in a good agreement with both the theoretical and experimental published data [3,10]. On the otherhand, the intermediate and thermal neutron fluxes show a maximum at around 2.5 cm inside the phantom. These maxima are attributed to the scattering processes occur inside the phantom between the diffused neutrons and hydrogen nuclei.

Making use of equation(2) and ICRP data [6], the absorbed doses and dose equivalent rates were determined and calculated for the different energy regions. These results are reproduced in figures (4,5). A good agreement was found between the determined and calculated doses. Also the variation of the dose build-up for integrated fast neutron with depth in the considered medium is represented in figure(6). It is clear that the dose build-up is dependent on the average scattering and the total macroscopic cross-sections. The data presented in this work isof great help in the field of neutron therapy.

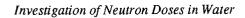
### ACKNOWLEDGEMENT

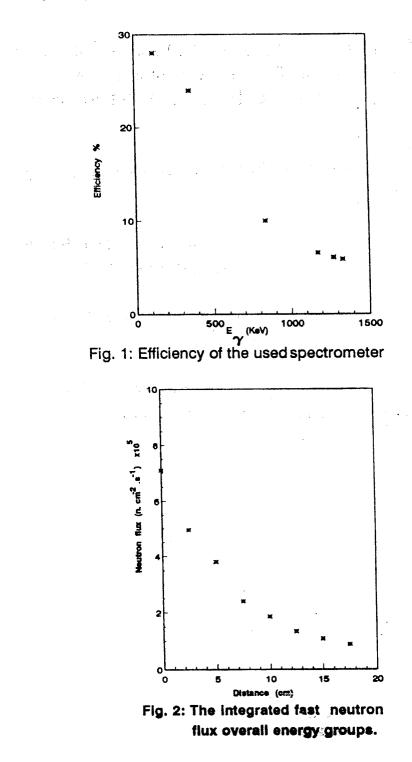
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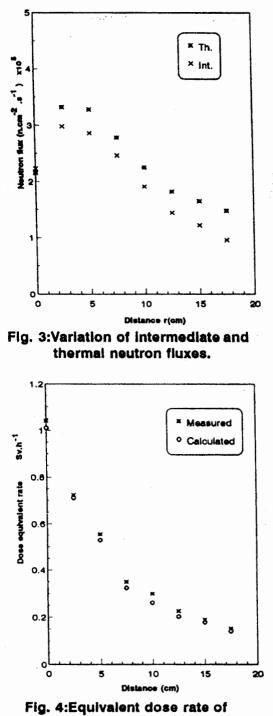
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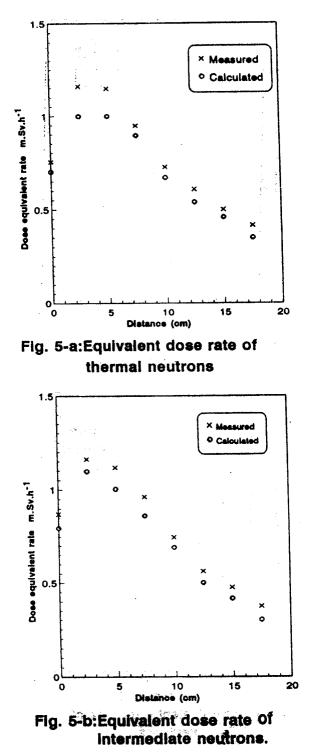




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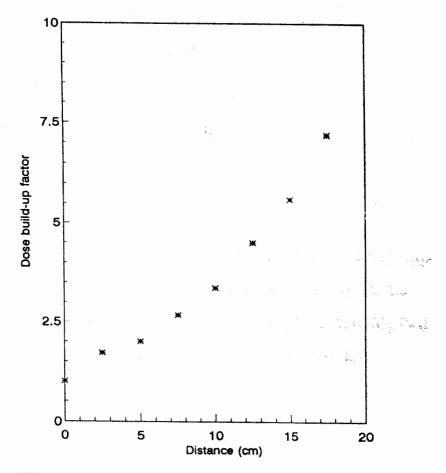


Fig. 6: Dose build-up for fast neutrons

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تحديد الجرعات النيوترونية داخل الماء محمد بسيوني قسم العلوم الأساسية – التطبيقية كلية الهندسة والتكنولوجيا – الأكاديمية العربية للعولوم والتكنولوجيا

# ملخص البحث

تم في هذا البحث تعيين التوزيع الطاقي والمكاني للسيل النيتروني المتولد من أنبوبة نيوترونية د + ت. وذلك عند انتشارها داخل جسم مكافيء للإنسان ( الماء) . وقد تمت القياسات بإستخدام طريقة التنشيط الأشعاعي للكواشف العتبية. حيث تم أستنتاج السيل النيوتروني المنتشر في داخل الجسم ومن ثم حسبت الجرعة الأشعاعية الكلية .